## Solution

## Week 51 (9/1/03)

## Accelerating spaceship

We will solve this problem by considering two nearby times and using the velocityaddition formula,

$$v = \frac{v_1 + v_2}{1 + v_1 v_2 / c^2} \,. \tag{1}$$

Using the definition of the proper acceleration, a, we have (with  $v_1 \equiv v(t)$  and  $v_2 \equiv a dt$ )

$$v(t+dt) = \frac{v(t) + a \, dt}{1 + v(t)a \, dt/c^2} \,. \tag{2}$$

Expanding both sides to first order in dt yields<sup>1</sup>

$$\frac{dv}{dt} = a\left(1 - \frac{v^2}{c^2}\right).\tag{3}$$

Separating variables and integrating gives, using  $1/(1-z^2) = 1/2(1-z) + 1/2(1+z)$ ,

$$\int_0^v \left(\frac{1}{1 - v/c} + \frac{1}{1 + v/c}\right) dv = 2a \int_0^t dt.$$
 (4)

This yields  $\ln \left( (1 + v/c)/(1 - v/c) \right) = 2at/c$ . Exponentiating, and solving for v, gives

$$v(t) = c \left(\frac{e^{2at/c} - 1}{e^{2at/c} + 1}\right) = c \tanh(at/c).$$

$$(5)$$

Note that for small a or small t (more precisely, for  $at/c \ll 1$ ), we obtain  $v(t) \approx at$ , as we should. And for  $at/c \gg 1$ , we obtain  $v(t) \approx c$ , as we should.

REMARKS: If a happens to be a function of time, a(t), then we can't move the *a* outside the integral in eq. (4), so we instead end up with the general formula,

$$v(t) = c \tanh\left(\frac{1}{c} \int_0^t a(t) \, dt\right). \tag{6}$$

If we define the *rapidity*,  $\phi$ , by

$$\phi(t) \equiv \frac{1}{c} \int_0^t a(t) \, dt,\tag{7}$$

then we have

$$v = c \tanh \phi \quad \Longleftrightarrow \quad \tanh \phi = \frac{v}{c} \,.$$
 (8)

Note that whereas v has c as a limiting value,  $\phi$  can become arbitrarily large. The  $\phi$  associated with a given v is simply 1/mc times the time integral of the force (felt by the astronaut) needed to bring the astronaut up to speed v. By applying a force for an arbitrarily long time, we can make  $\phi$  arbitrarily large.

<sup>&</sup>lt;sup>1</sup>Equivalently, just take the derivative of  $(v + w)/(1 + vw/c^2)$  with respect to w, and then set w = 0.

The quantity  $\phi$  is very useful because many expressions in relativity (which we'll just invoke here) take on a particularly nice form when written in terms of  $\phi$ . Consider, for example, the velocity-addition formula. Let  $\beta_1 = \tanh \phi_1$  and  $\beta_2 = \tanh \phi_2$ . Then if we add  $\beta_1$  and  $\beta_2$  using the velocity-addition formula, eq. (1), we obtain

$$\frac{\beta_1 + \beta_2}{1 + \beta_1 \beta_2} = \frac{\tanh \phi_1 + \tanh \phi_2}{1 + \tanh \phi_1 \tanh \phi_2} = \tanh(\phi_1 + \phi_2),\tag{9}$$

where we have used the addition formula for  $\tanh \phi$  (which can be proved by writing things in terms of the exponentials  $e^{\pm \phi}$ ). Therefore, while the velocities add in the strange manner of eq. (1), the rapidities add by standard addition.

The Lorentz transformation,

$$\begin{pmatrix} x \\ ct \end{pmatrix} = \begin{pmatrix} \gamma & \gamma\beta \\ \gamma\beta & \gamma \end{pmatrix} \begin{pmatrix} x' \\ ct' \end{pmatrix},$$
(10)

also takes a nice form when written in terms of the rapidity. Note that  $\gamma$  can be written as

$$\gamma \equiv \frac{1}{\sqrt{1-\beta^2}} = \frac{1}{\sqrt{1-\tanh^2 \phi}} = \cosh \phi, \tag{11}$$

and so

$$\gamma\beta \equiv \frac{\beta}{\sqrt{1-\beta^2}} = \frac{\tanh\phi}{\sqrt{1-\tanh^2\phi}} = \sinh\phi.$$
(12)

Therefore, the Lorentz transformation becomes

$$\begin{pmatrix} x \\ ct \end{pmatrix} = \begin{pmatrix} \cosh \phi & \sinh \phi \\ \sinh \phi & \cosh \phi \end{pmatrix} \begin{pmatrix} x' \\ ct' \end{pmatrix}.$$
 (13)

This looks similar to a rotation in a plane, which is given by

$$\begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{pmatrix} \begin{pmatrix} x' \\ y' \end{pmatrix},$$
(14)

except that we now have hyperbolic trig functions instead of the usual trig functions. The fact that the *invariant interval*,  $s^2 \equiv c^2 t^2 - x^2$ , does not depend on the frame is clear from eq. (13), because the cross terms in the squares cancel, and  $\cosh^2 \phi - \sinh^2 \phi = 1$ . (Compare with the invariance of  $r^2 \equiv x^2 + y^2$  for rotations in a plane.)

Quantities associated with a Minkowski diagram also take a nice form when written in terms of the rapidity. In particular, the angle between the axes of the two relevant frames happens to be  $\tan \theta = \beta$ , where  $\beta c$  is the relative speed between the frames. But  $\beta = \tanh \phi$ , so the angle between the axes is given by

$$\tan \theta = \tanh \phi. \tag{15}$$

The integral  $\int a(t) dt$  (which is c times the rapidity) may be described as the naive, incorrect speed. That is, it is the speed the astronaut might *think* he has, if he has his eyes closed and knows nothing about the theory of relativity. (And indeed, his thinking would be essentially correct for small speeds.) The quantity  $\int a(t) dt$  seems like a reasonably physical thing, so if there is any justice in the world,  $\int a(t) dt = \int F(t) dt/m$  should have *some* meaning. And indeed, although it doesn't equal v, all you have to do to get v is take a tanh and throw in some factors of c.

The fact that rapidities add via simple addition when using the velocity-addition formula, as we saw in eq. (9), is evident from eq. (6). There is really nothing more going on here than the fact that

$$\int_{t_0}^{t_2} a(t) \, dt = \int_{t_0}^{t_1} a(t) \, dt + \int_{t_1}^{t_2} a(t) \, dt. \tag{16}$$

To be explicit, let a force be applied from  $t_0$  to  $t_1$  that brings a mass up to speed  $\beta_1 = \tanh \phi_1 = \tanh(\int_{t_0}^{t_1} a \, dt)$ , and then let an additional force be applied from  $t_1$  to  $t_2$  that adds on an additional speed of  $\beta_2 = \tanh \phi_2 = \tanh(\int_{t_1}^{t_2} a \, dt)$  (relative to the speed at  $t_1$ ). Then the resulting speed may be looked at in two ways: (1) it is the result of relativistically adding the speeds  $\beta_1 = \tanh \phi_1$  and  $\beta_2 = \tanh \phi_2$ , and (2) it is the result of applying the force from  $t_0$  to  $t_2$  (you get the same final speed, of course, whether or not you bother to record the speed along the way at  $t_1$ ), which is  $\beta = \tanh(\int_{t_0}^{t_2} a dt) = \tanh(\phi_1 + \phi_2)$ , where the last equality comes from the obvious statement, eq. (16). Therefore, the relativistic addition of  $\tanh \phi_1$  and  $\tanh \phi_2$  gives  $\tanh(\phi_1 + \phi_2)$ , as we wanted to show.