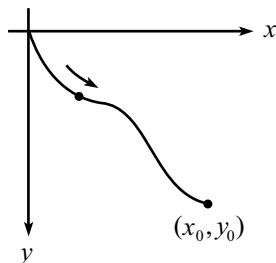


Solution

Week 83 (4/12/04)

The brachistochrone

First solution: In the figure below, the boundary conditions are $y(0) = 0$ and $y(x_0) = y_0$, with downward taken to be the positive y direction.



From conservation of energy, the speed as a function of y is $v = \sqrt{2gy}$. The total time is therefore

$$T = \int_0^{x_0} \frac{ds}{v} = \int_0^{x_0} \frac{\sqrt{1+y'^2}}{\sqrt{2gy}} dx. \quad (1)$$

Our goal is to find the function $y(x)$ that minimizes this integral, subject to the boundary conditions above. We can therefore apply the results of the variational technique, with a “Lagrangian” equal to

$$L \propto \frac{\sqrt{1+y'^2}}{\sqrt{y}}. \quad (2)$$

The Euler-Lagrange equation is

$$\frac{d}{dx} \left(\frac{\partial L}{\partial y'} \right) = \frac{\partial L}{\partial y} \quad \implies \quad \frac{d}{dx} \left(\frac{1}{\sqrt{y}} \cdot y' \cdot \frac{1}{\sqrt{1+y'^2}} \right) = -\frac{\sqrt{1+y'^2}}{2y\sqrt{y}}. \quad (3)$$

Using the product rule on the three factors on the left-hand side, and making copious use of the chain rule, we obtain

$$-\frac{y'^2}{2y\sqrt{y}\sqrt{1+y'^2}} + \frac{y''}{\sqrt{y}\sqrt{1+y'^2}} - \frac{y'^2 y''}{\sqrt{y}(1+y'^2)^{3/2}} = -\frac{\sqrt{1+y'^2}}{2y\sqrt{y}}. \quad (4)$$

Multiplying through by $2y\sqrt{y}(1+y'^2)^{3/2}$ and simplifying gives

$$-2yy'' = 1 + y'^2. \quad (5)$$

We can integrate this equation if we multiply through by y' and rearrange to obtain

$$\begin{aligned} \int \frac{2y'y''}{1+y'^2} &= - \int \frac{y'}{y} \cdot & \implies & \ln(1+y'^2) = -\ln y + A \\ & & \implies & 1+y'^2 = \frac{B}{y}, \end{aligned} \quad (6)$$

where $B \equiv e^A$. We must now integrate one more time. Solving for y' and separating variables gives

$$\frac{\sqrt{y} dy}{\sqrt{B-y}} = \pm dx. \quad (7)$$

A helpful change of variables to get rid of the square root in the denominator is $y \equiv B \sin^2 \phi$. Then $dy = 2B \sin \phi \cos \phi d\phi$, and eq. (7) simplifies to

$$2B \sin^2 \phi d\phi = \pm dx. \quad (8)$$

We can now make use of the relation $\sin^2 \phi = (1 - \cos 2\phi)/2$ to integrate this. The result is $B(2\phi - \sin 2\phi) = \pm 2x - C$, where C is an integration constant.

Now note that we may rewrite our definition of ϕ (which was $y \equiv B \sin^2 \phi$) as $2y = B(1 - \cos 2\phi)$. If we then define $\theta \equiv 2\phi$, we have

$$x = \pm a(\theta - \sin \theta) \pm d, \quad y = a(1 - \cos \theta). \quad (9)$$

where $a \equiv B/2$, and $d \equiv C/2$.

The particle starts at $(x, y) = (0, 0)$. Therefore, θ starts at $\theta = 0$, since this corresponds to $y = 0$. The starting condition $x = 0$ then implies that $d = 0$. Also, we are assuming that the wire heads down to the right, so we choose the positive sign in the expression for x . Therefore, we finally have

$$x = a(\theta - \sin \theta), \quad y = a(1 - \cos \theta). \quad (10)$$

This is the parametrization of a *cycloid*, which is the path taken by a point on the rim of a rolling wheel. The initial slope of the $y(x)$ curve is infinite, as you can check.

Second solution: Let's use a variational argument again, but now with y as the independent variable. That is, let the chain be described by the function $x(y)$. The arclength is now given by $ds = \sqrt{1 + x'^2} dy$. Therefore, instead of the Lagrangian in eq. (2), we now have

$$L \propto \frac{\sqrt{1 + x'^2}}{\sqrt{y}}. \quad (11)$$

The Euler-Lagrange equation is

$$\frac{d}{dy} \left(\frac{\partial L}{\partial x'} \right) = \frac{\partial L}{\partial x} \quad \Longrightarrow \quad \frac{d}{dy} \left(\frac{1}{\sqrt{y}} \frac{x'}{\sqrt{1 + x'^2}} \right) = 0. \quad (12)$$

The zero on the right-hand side makes things nice and easy, because it means that the quantity in parentheses is a constant. Call it D . We then have

$$\begin{aligned} \frac{1}{\sqrt{y}} \frac{x'}{\sqrt{1 + x'^2}} = D &\quad \Longrightarrow \quad \frac{1}{\sqrt{y}} \frac{dx/dy}{\sqrt{1 + (dx/dy)^2}} = D \\ &\quad \Longrightarrow \quad \frac{1}{\sqrt{y}} \frac{1}{\sqrt{(dy/dx)^2 + 1}} = D. \end{aligned} \quad (13)$$

This is equivalent to eq. (6), and the solution proceeds as above.

Third solution: Note that the “Lagrangian” in the first solution above, which is given in eq. (2) as

$$L = \frac{\sqrt{1 + y'^2}}{\sqrt{y}}, \quad (14)$$

is independent of x . Therefore, in analogy with conservation of energy (which arises from a Lagrangian that is independent of t), the quantity

$$E \equiv y' \frac{\partial L}{\partial y'} - L = \frac{y'^2}{\sqrt{y}\sqrt{1 + y'^2}} - \frac{\sqrt{1 + y'^2}}{\sqrt{y}} = \frac{-1}{\sqrt{y}\sqrt{1 + y'^2}} \quad (15)$$

is a constant. We have therefore again reproduced eq. (6).